

STRONG PERTURBATIONS ON THE SUN AND IN THE HELIOSPHERE: SCALING OF SIMILAR AND INDIVIDUAL CHARACTERISTICS

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Contents

- Dimensionless parameters in MHD with dissipation and radiation
- Eruptions with/without topological changes ('reconnection')
- Velocity scaling and velocity 'reconnection'
- Classical case interpretation (eruption of May 15, 1919 on the Sun)
- Conclusions

Dimensionless scaling: Alfvén wave

- **Given:**

- 1) magnetic field B ;

- 2) mass density ρ

- **Problem:**

Find characteristic velocity of perturbations

- **Answer:**

$$V_A \sim B/\rho^{1/2}$$

(H. Alfvén, 1942)

Dimensionless scaling: strong explosions in the atmosphere (Sedov's solution)

- **Given:**

- 1) Strong point-like explosion in the homogeneous atmosphere
- 2) E - energy released
- 3) ρ - density of the atmosphere
- 4) t – time after explosion

- **Problem:**

Find the shock velocity $u(t)$ as a function of time

- **Solution:**

Dimensionless parameter was found (L.I. Sedov) $\xi = r (\rho/E t^2)^{1/5}$

- **Answer:**

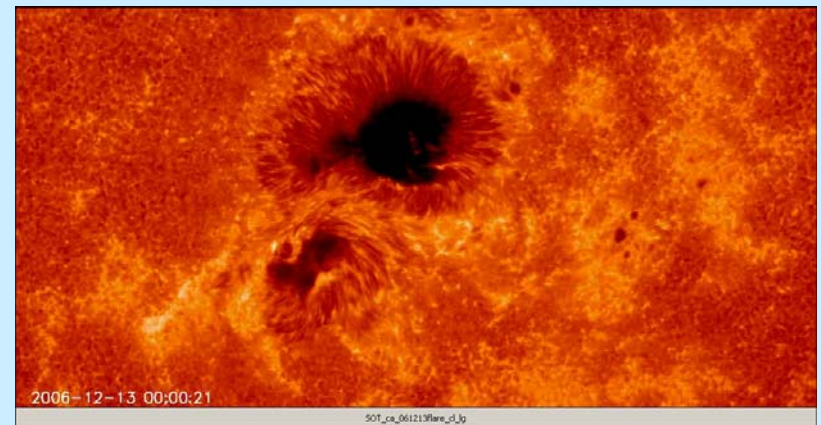
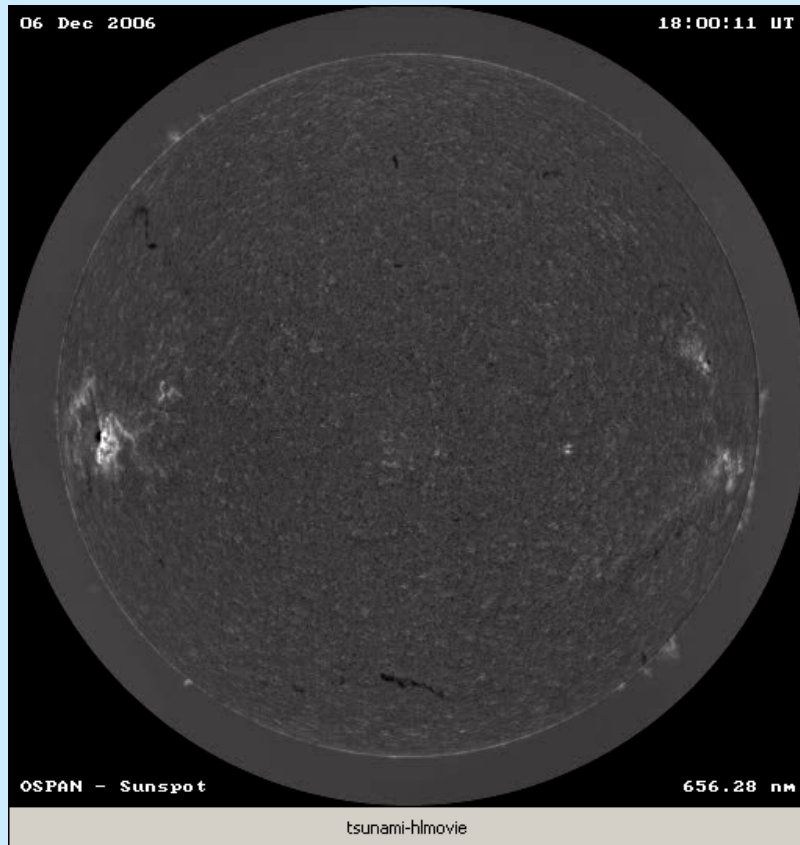
$$u(t) \sim (E/\rho t^3)^{1/5}$$

- **Practical implications: Energy of the explosion E can be estimated from kinematics of the strong shocks**

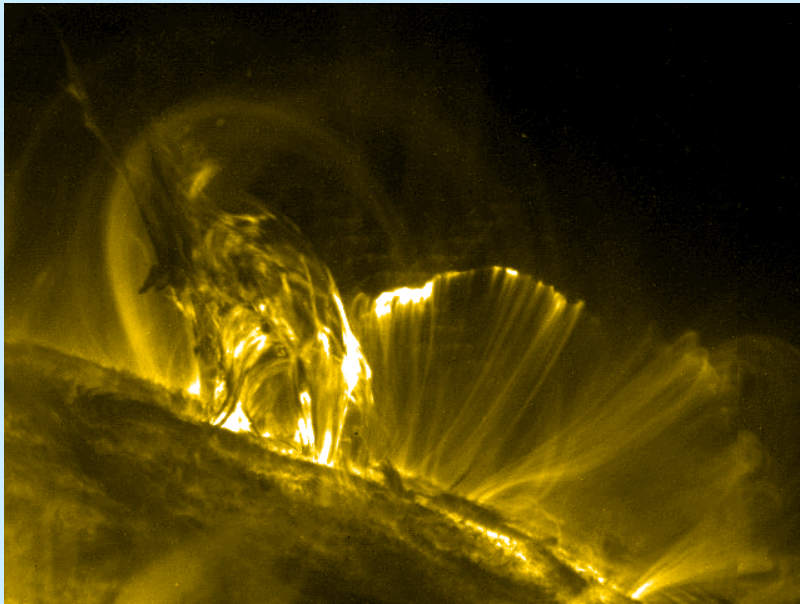
Explosive wave

Hinode observations

December 6, 2006 (656.28 nm)



Complicated dynamical geometry of arcades:



- Spurious “reconnections” and X-points: projections of twisted loops
- Spurious “crossings of the loops”: the same as above
- Spurious “plasmoids”: brighter summits of the loops

(an example from the TRACE “gold mine”)

Projection effects



Multi-scale nature of erupting prominences



Conservation Laws

mass

momentum

energy

Governing MHD equations with dissipation and radiation

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0, \quad (1)$$

$$\begin{aligned} \rho \left\{ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \vec{\nabla}) \vec{v} \right\} + \vec{\nabla} p - \frac{1}{4\pi} [[\vec{\nabla} \times \vec{B}] \times \vec{B}] + \\ + \vec{F}_{viscous} + \vec{F}_{gravity} = 0 \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{\partial}{\partial t} \left(\rho u + \frac{1}{2} \rho v^2 \right) + \frac{\partial}{\partial x_i} \left(\rho v_i w + \rho v_i \frac{v^2}{2} + \rho v_k \sigma'_{ik} - \kappa \frac{\partial T}{\partial x_i} + \right. \\ \left. + \frac{1}{4\pi} e_{ikl} e_{lmn} v_m B_k B_n + \frac{c}{4\pi\sigma} e_{ikl} j_k B_l \right) = -L + A \end{aligned} \quad (3)$$

Symbolic shape of equations

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0, \quad (4)$$

$$\begin{aligned} S^{-1} + 1 \\ \rho \left\{ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right\} = - \vec{\nabla} p + \frac{1}{4\pi} [[\vec{\nabla} \times \vec{B}] \times \vec{B}] + \\ + \vec{F}_{viscous} + \vec{F}_{gravity} \end{aligned} \quad (5)$$

$$S^{-1} + 1 = M^{-2} + M_A^{-2} + Re^{-1} + Fr^{-1}$$

$$\begin{aligned} \frac{\partial}{\partial t} (\rho u + \frac{1}{2} \rho v^2) + \frac{\partial}{\partial x_i} (\rho v_i w + \rho v_i \frac{v^2}{2} + \rho v_k \sigma'_{ik} - \kappa \frac{\partial T}{\partial x_i} + \\ + \frac{1}{4\pi} e_{ikl} e_{lmn} v_m B_k B_n + \frac{c}{4\pi \sigma} e_{ikl} j_k B_l) = -L + A \end{aligned} \quad (6)$$

$$\begin{aligned} M^{-2} S^{-1} + S^{-1} + M^{-2} + 1 + Re^{-1} + Re^{-1} M^{-2} Pr^{-1} \\ + M_A^{-1} + M_A^{-2} Re_m^{-1} = V_e + Fr^{-1} \end{aligned}$$

Dimensionless parameters

Name	Description	Role
Strouhal	Time / Flight times	Time scales
Knudsen	Mean free path / Size	Length scales
Velocity-emission	Kinetic energy / EM emission	Plasma density
Mach	Bulk speed / Thermal speed	Temperature
Magnetic Mach	Bulk speed / Alfvén speed	Magnetic field
Froude	Velocity / Free escape speed	Gravity
Faraday	Potential / Inductive	Electric field
Trieste numbers	Inflows / Inner flows	Openness degrees

Table 3. Irreducible and full set of independent orthogonal dimensionless parameters in the ‘velocity normalised basis’

Symbol	Name	Formula / Definition	Significance
M	Mach number	v/c_s	Temperature
Ma	Magnetic Mach (Mach – Alfvén)	v/V_A	Magnetic field and currents
S	Strouhal	vt/x	Time
Tr	Trieste numbers (set)	Ratio of internal and boundary crossing fluxes	Openness degree of the object against energy, momentum and mass fluxes
Ve	“Velocity-emission”	$\left(\frac{E_{kin}}{E_{at}}\right)\left(\frac{x}{v\tau_{rad}}\right)$	Mass density
Kn	Knuden numbers (set)	λ/x	Length
Fr	Froude	v^2/gl	Gravity
F	Faraday number	$\frac{1}{c} \frac{\partial A}{\partial t} / \nabla \varphi$	Electric fields and charges $\left(\frac{j}{\rho c}\right)\left(\frac{x}{ct}\right)$

Table 2. Binary classification of proper and external flow types

Types/ motions	SS	SF	FS	FF
Proper motions	Slow	Slow	Fast	Fast
External motions	Slow	Fast	Slow	Fast

New interpretation of the classical observation

**E. Pettit, The Great Eruptive Prominences of May
29 and July 15, 1919,
Astrophys. J., vol. 50, pp. 206-219, 1919**

Internal and External Motions

Velocity reconnection

Loops and whips

E.Pettit does not admitted external motions

May 15, 1919

huge prominence on the East limb

8 hours rise up to 2 solar radii

stereo couples and velocity tracing

a loop transformed into a whip

Figure 2

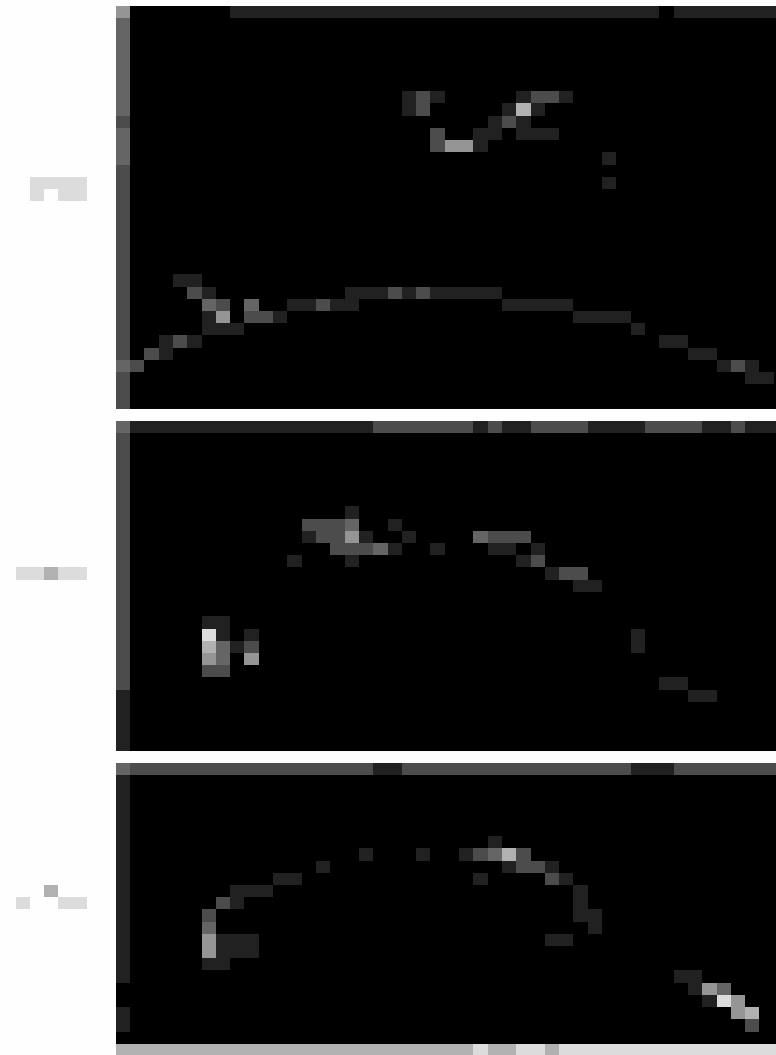


Figure 2: Evolution of a wave packet over time.

Height versus time



Detailed stereoscopic velocity measurements (0 - ~ 100 km/s)



Simple velocity model
siphon flow+solar wind=
whip formation

Strouhal number

- $S \sim vt/L$ **non-invariant, reference frame dependent parameter**
- $S \gg 1$ quasi-stationary flow (asymptotically large S values mean steady state)
- $S \ll 1$ transient flow
- $S \sim 1$ intermediate regime (for example, linear waves or small convective perturbations)

Strouhal number (ctnd)

- $S \sim Vt/L$
- Quiescent prominences:
- $L \sim 10^5$ km, $V \sim 1$ km/s, $t \sim 10^6$ s (many days, a month) (life time)
- Hence, $S \sim 10$
- Looks like quasi-steady flows during life time

Strouhal number (ctnd)

- $S \sim vt/L$
- Eruptive prominences: for example,
- $L \sim 10^5$ km, $V \sim 100$ km/s, $t \sim 1$ hour
- Looks like transient phenomenon

Laminar? Turbulent?

- Number of degrees of freedom N involved in process under consideration
- Laminar if $N \sim 1$ “simple”
- Turbulent if $N \gg 1$ “complicated”
- Reynolds number Re is not invariant for this purpose (reference frame dependent impressions of flow). Because of this,
- $Re > 1$, $Re < 1$ is not a good classification.

Dimensionless Faraday number F

- Ratio of inductive and potential electric fields
- $F \sim (j/\rho c) (L/ct)$

where

j - electric current density

ρ - electric charge density

L - length

c - velocity of light

t - time

$F \gg 1$ inductive fields dominate

$F \ll 1$ Coulomb fields dominate \rightarrow electrostatics **important!**

Applications: thermonuclear fusion problem, double layers
on the Sun

Trieste numbers T

- Set of dimensionless parameters: ratios of energy, momentum, mass fluxes inside, outside and across the volume boundaries
- $T \gg 1$ open system
- $T \ll 1$ closed system
- $T \sim 1$ intermediate case
- **Important example:** quiescent prominence – open system with mass flows through it (siphon flows in loops from one foot point to opposite foot point, blue and red shifts in legs). It is not a magneto-static equilibrium (as generally believed) but a quasi-steady or non-steady state with flows

Why solar flares and CMEs originate?

- It is because of subphotospheric free energy supply in the same, larger and smaller space-time scales: $T \sim 1$, $T \gg 1$, $T \ll 1$.
- All these regimes involved in preparation and development
- They are non-linearly coupled
- Information and ‘signals’ from interior of the Sun are needed to predict solar flares and CMEs

V_e number (on the Sun)

- Ratio of the energy losses via solar wind to the losses due to electromagnetic radiation of the corona
- Averages on time :
- $V_e \gg 1$ in coronal holes (dark regions, fast wind)
- $V_e \ll 1$ in active regions (bright loops, slow motions)
- $V_e \sim 1$ intermediate value in quiet corona

V_e number (continued)

- $V_e \gg 1$ in coronal mass ejections (CMEs) – high loop eruptions, prominences, arcades
- $V_e \ll 1$ in solar flares – low loops, brightening, confined motions
- $V_e \sim 1$ intermediate case - both
- Typically CMEs and flares are accompanying each other in different proportions from CME-like to flare like situations
- Flares/CMEs: No cause/consequence relations, but parallel energy release channels from the same free energy sources

Syrovatsky number

- Defined as combination M^2/M_A .
- Includes temperature, magnetic field, velocity
- NOT INCLUDES DENSITY!

Comfortable for the classification of incompressible MHD flows

- The same $M \beta^{-1/2}$
- The same $M_A \beta^{-1}$
- Here $\beta \sim (M_A/M)^2 \cdot P/B^2$

CONCLUSIONS

- ✓ Dimensionless velocity scaling gives unambiguous classification of ejection on the Sun.
- ✓ 8 binary and 27 ternary types of the velocity field can be indicated.
- ✓ Velocity field reconnection is observed during ejection as “loop - whip shape transformation.”
- ✓ Classical observations by E. Pettit are interpreted using this approach: siphon+solar wind= whip.
- ✓ **Thank you!**